

Erratum: “The Band Gap Formation in Rotors With Longitudinal Periodicity and Quasi-Periodicity” [ASME J. Eng. Gas Turbines Power, 2022, 144(5), p. 051003; DOI: [10.1115/1.4053193](https://doi.org/10.1115/1.4053193)]

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Equation (10) gives

$$\sum_{j=1}^n f(x_j) \approx \frac{1}{\Delta x} \int_0^L f(x) dx$$

where

$$f(x_j) = \frac{m_j}{\rho AL} \phi_i(x_j) \phi_k(x_j)$$

Considering that all disks have the same mass m_D , then

$$f(x_j) = \frac{m_D}{\rho AL} \phi_i(x_j) \phi_k(x_j)$$

Hence, from Eq. (10)

$$\begin{aligned} \sum_{j=1}^{N_D} \frac{m_D}{\rho AL} \phi_i(x_j) \phi_k(x_j) &\approx \frac{1}{\Delta x} \int_0^L \frac{m_D}{\rho AL} \phi_i(x) \phi_k(x) dx \\ &= \frac{\hat{m}}{\Delta x} L \delta_{ik} = \hat{m}(N_D + 1) \delta_{ik} = \mu \delta_{ik} \end{aligned} \quad (11)$$

where $\hat{m} = m_D/\rho AL$ and $\mu = \hat{m}(N_D + 1)$, which we can call the specific disk mass.

Similarly, if we consider that all disks have the same inertia J_D , then

$$\sum_{j=1}^{N_D} \frac{J_D}{\rho AL} \phi_i(x_j) \frac{d^2 \phi_k(x_j)}{dx^2} \approx \frac{1}{\Delta x} \int_0^L \frac{J_D}{\rho AL} \phi_i(x) \frac{d^2 \phi_k(x)}{dx^2} dx \quad (12)$$

Therefore, considering Eqs. (13) and (14)

$$\begin{aligned} \sum_{j=1}^{N_D} \frac{J_D}{\rho AL} \phi_i(x_j) \phi_k(x_j) &\approx \frac{1}{\Delta x} \int_0^L -\frac{J_D}{\rho AL} \left(\frac{k\pi}{L}\right)^2 \phi_i(x) \phi_k(x) dx \\ &= -\hat{J}(N_D + 1) \left(\frac{k\pi}{L}\right)^2 \delta_{ik} = \lambda \left(\frac{k\pi}{L}\right)^2 \delta_{ik} \end{aligned} \quad (15)$$

where $\hat{J} = J_D/\rho AL$ and $\lambda = \hat{J}(N_D + 1)$, which we can call the specific disk inertia.

The results shown in Figs. 15–17 were obtained using the present formulation of μ and λ .